

Fig. 1. Lower part: measured slopes  $\partial \ln T_c / \partial p$ . For  $p > 10$  kbar comparison is made with  $\partial \ln \tau / \partial n$ . Relevant parameters are plotted in the upper part

In the Table 1 we list the relevant new data. We note that  $T_c$  of our Ta is rather low.

The data on the effect of *pressure* on these alloys are shown in Fig. 1, including values for  $T_c$  and the Debye characteristic temperature  $\theta$ , if known<sup>8,9</sup>. For pressures in excess of 10 kbar, the slope  $1/T_{c0} \cdot \partial T_c / \partial p \approx \partial \ln T_c / \partial p$  decreases monotonically with increasing number  $n$  of valence electrons per atom. Proceeding as for the Zr-Nb-Mo alloys, we have plotted the quantity  $a \cdot \partial \ln \tau / \partial n \equiv a(\partial \ln T_c / \partial n - \partial \ln \theta / \partial n)$ .

<sup>9</sup> We wish to thank G. Dammer and K. Mulder of this Laboratory for permitting us to use their unpublished data for Hf<sub>61</sub>Ta<sub>39</sub>.

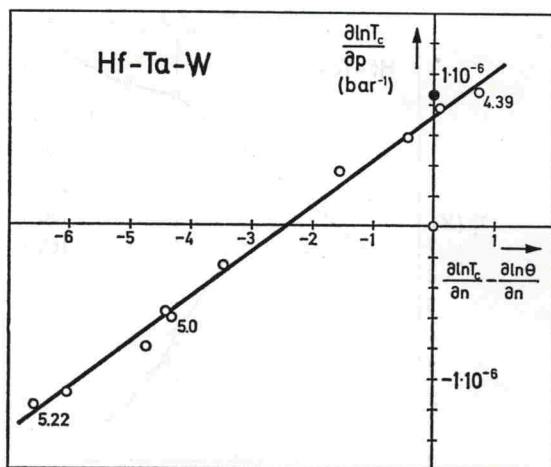


Fig. 2. Plot showing the linear relation between  $\partial \ln T_c / \partial p$  and  $\partial \ln \tau / \partial n$  for all alloys. The full circle represents  $\partial \ln \theta / \partial p$  for Ta

One notes that for Hf-Ta-W-alloys it is also linearly related with  $\partial \ln T_c / \partial p$ , according to

$$\partial \ln T_c / \partial p = a(\partial \ln T_c / \partial n - \partial \ln \theta / \partial n) + b. \quad (2)$$

This is also shown in Fig. 2, from which the values of the constants  $a = 0.30 \cdot 10^{-6} \text{ bar}^{-1}$  and  $b = 0.74 \cdot 10^{-6} \text{ bar}^{-1}$  are readily obtained.

Inserting  $T_c \propto \theta \exp(-1/g)$  in Eq. (2) yields

$$g^{-2} \cdot \partial g / \partial p + \partial \ln \theta / \partial p = g^{-2} \cdot a \cdot \partial g / \partial n + b. \quad (3)$$

For Ta the literature value of  $\partial \ln \theta / \partial p = \gamma_G \cdot \kappa = 1.82 \cdot 0.48 \cdot 10^{-6} \text{ bar}^{-1} = 0.87 \cdot 10^{-6} \text{ bar}^{-1}$  (10) is close to the value of our term  $b = 0.74 \cdot 10^{-6} \text{ bar}^{-1}$  ( $\gamma_G$  = Grüneisen parameter,  $\kappa$  = compressibility). Thus, under the reasonable assumption that  $\partial \ln \theta / \partial p$  does not change much with composition, Eq. (3) reduces to Eq. (1), with  $a = (0.30 \pm 0.02) \cdot 10^{-6} \text{ bar}^{-1}$ .

For Ta-W alloys kinks in the otherwise linear  $T_c(p)$ -dependence are observed near  $p = 10$  kbar, as indicated by the splitting of  $\partial \ln T_c / \partial p$  in Fig. 1. A similar behaviour was observed in the Zr-Nb-Mo system. However, Nb<sub>96</sub>Zr<sub>4</sub> and Nb were also involved and the changes in slope occurred near 20 kbar. We will refer to this effect in another paper.

10 Gschneidner, K. A.: In: Solid State Physics 16, 275 (1964), edited by F. Seitz and D. Turnbull. London-New York: Academic Press.